

Solution and Marking Scheme

Theory

I. Satellite's orbit transfer

a) $\frac{mu_0^2}{R_0} = \frac{GMm}{R_0^2}, \quad u_0 = \sqrt{\frac{GM}{R_0}}$ (1 point)

b) conservation of angular momentum: $mu_1R_0 = mu_2R_1$
 conservation of energy: $\frac{1}{2}mu_2^2 - \frac{GMm}{R_1} = \frac{1}{2}mu_1^2 - \frac{GMm}{R_0}$

$$\left[\left(\frac{R_0}{R_1} \right)^2 - 1 \right] u_1^2 = 2GM \left[\frac{1}{R_1} - \frac{1}{R_0} \right]$$

$$\frac{(R_0 - R_1)(R_0 + R_1)}{R_1^2} u_1^2 = (2GM) \frac{(R_0 - R_1)}{R_0 R_1}$$

$$u_1 = \sqrt{\frac{GM}{R_0}} \sqrt{\frac{2R_1}{R_1 + R_0}} = u_0 \sqrt{\frac{2R_1}{R_1 + R_0}} \quad (2 \text{ points})$$

c) $\lim_{R_1 \rightarrow \infty} u_1 = \sqrt{2} u_0$ (1 point)

d) $u_2 = u_1 \frac{R_0}{R_1} = u_0 \frac{\sqrt{2} R_0}{\sqrt{R_1(R_1 + R_0)}}$ (1 point)

e) $u_3 = \sqrt{\frac{GM}{R_1}} = \sqrt{\frac{GM}{R_0}} \sqrt{\frac{R_0}{R_1}} = u_0 \sqrt{\frac{R_0}{R_1}}$
 $= \sqrt{\frac{R_0}{R_1}} \sqrt{\frac{R_1(R_1 + R_0)}{\sqrt{2} R_0}} u_2$
 $u_3 = u_2 \sqrt{\frac{R_1 + R_0}{2R_0}} \quad (1 \text{ point})$

f) (3 points) combining equations (1) and (2) :

$$\frac{d^2}{dt^2} r - \frac{C/m}{r^3} = -\frac{GM}{r^2}$$

and for the circular orbit of radius R_1 we have $\frac{C}{m} = GMR_1$

hence $\frac{d^2}{dt^2} r - \frac{GMR_1}{r^3} = -\frac{GM}{r^2}$

putting $r = R_1 + \eta$, where $\eta \ll R_1$

$$\therefore \frac{d^2}{dt^2} \eta - \frac{GMR_1}{R_1^3 \left(1 + \frac{\eta}{R_1} \right)^3} = -\frac{GM}{R_1^2 \left(1 + \frac{\eta}{R_1} \right)^2}$$

$$\frac{d^2}{dt^2}\eta - \frac{GM}{R_1^2} \left(1 - 3\frac{\eta}{R_1}\right) \approx -\frac{GM}{R_1^2} \left(1 - 2\frac{\eta}{R_1}\right)$$

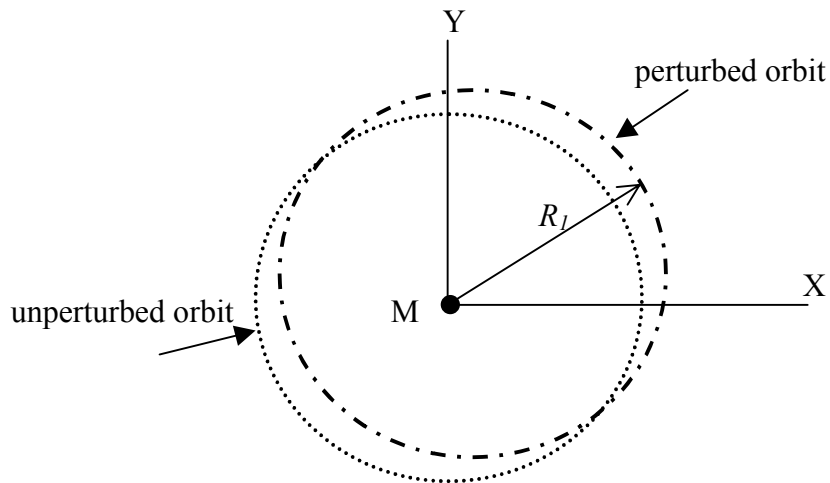
$$\frac{d^2}{dt^2}\eta \approx -\frac{GM}{R_1^3}\eta$$

the frequency of oscillation about mean distance is $f = \frac{1}{2\pi} \sqrt{\frac{GM}{R_1^3}}$

the period $T = \frac{1}{f} = 2\pi \sqrt{\frac{R_1^3}{GM}}$

Note that this period is the same as the orbital period

h) (1 point)



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II. Optical Gyroscope

The light wave moves with speed $c' = \frac{c}{\mu}$ in the medium having refractive index μ . Wavelength of light in medium $\lambda' = \frac{\lambda}{\mu}$, where λ is the wavelength of light in vacuum.

a) (2 points)

$$\text{transit time for the CW beam: } t^+ = \frac{2\pi R + R\Omega t^+}{c'} = \frac{2\pi R}{c'} \left(1 - \frac{R\Omega}{c'}\right)^{-1}$$

$$\text{transit time for the CCW beam: } t^- = \frac{2\pi R - R\Omega t^-}{c'} = \frac{2\pi R}{c'} \left(1 + \frac{R\Omega}{c'}\right)^{-1}$$

$$\text{the time difference between } t^+ \text{ and } t^-: \Delta t = \frac{4\pi R^2 \Omega}{(c')^2 - R^2 \Omega^2}$$

$$\text{since } (R\Omega)^2 \ll (c')^2 \quad \Delta t \approx \frac{4\pi R^2 \Omega}{(c')^2}$$

b) (2 points) the round-trip optical path difference, ΔL , is given by

$$\Delta L = c' \Delta t = \frac{4\pi R^2 \Omega}{c'}$$

c) (1 point) $\Delta L \cong 4.5 \times 10^{-12} \text{ m}$.

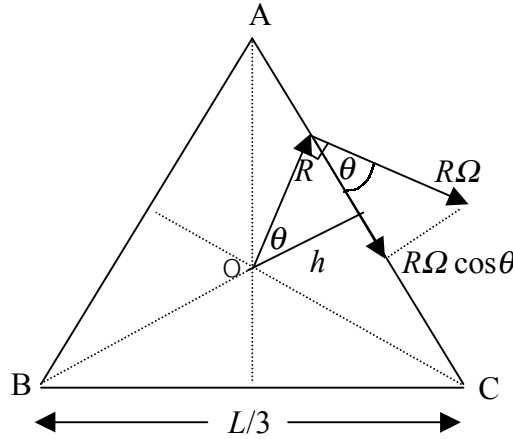
d) (1 point) the corresponding optical phase difference $\Delta\theta$ is,

$$\Delta\theta = \frac{2\pi \Delta L}{\lambda'} = \frac{8\pi^2 R^2 \Omega}{c\lambda'}, \text{ where } \lambda' = \frac{\lambda}{\mu}$$

for N turns of fiber optic ring,

$$\Delta\theta = \frac{8\pi^2 R^2 N \Omega}{c\lambda'}$$

e) (2 points)



The figure shows the triangular ring rotating about the centre o with the angular speed Ω in the clockwise direction. Without loosing generality, let's first consider the velocity of light along AC in the CW and CCW direction,

$$v_{\pm} = c \pm R\Omega \cos \theta = c \pm \Omega h, \text{ where } h \text{ is constant.}$$

$$\tau_{\pm} = \frac{L/3}{v_{\pm}} = \frac{L/3}{c \pm \Omega h} \approx \frac{L/3}{c} \left(1 \mp \frac{\Omega h}{c}\right)$$

where $\tau_{\pm}$ is the time taken for light travelling along AC in the CW and CCW.

$$t_{\pm} = \frac{L}{v_{\pm}} = \frac{L}{c \pm \Omega h} \approx \frac{L}{c} \left(1 \mp \frac{\Omega h}{c}\right), \text{ where } L \text{ is the perimeter of the triangular}$$

ring.

Therefore, the time difference of light travelling in one complete cycle.

$$\Delta t = \frac{2\Omega L h}{c^2} = \frac{4\Omega}{c^2} \left(\frac{1}{2} L h\right) = \frac{4\Omega A}{c^2}, \text{ where } A \text{ is the area of the triangular ring.}$$

f) The resonance frequencies associated with $L_{\pm}$ corresponding to the effective cavity lengths seen by CW and CCW propagating beams respectively is,

$$L_{+} = ct^{+} \approx L \left(1 - \frac{\Omega h}{c}\right)$$

$$L_{-} = ct^{-} \approx L \left(1 + \frac{\Omega h}{c}\right)$$

where $L_{\pm}$ is the perimeter of the equilateral triangle in the CW (+) and CCW (-) and we also use the fact that $h\Omega \ll c$. Therefore,

$$\Delta L = L_{-} - L_{+} = 2L \frac{\Omega h}{c} = \frac{4\Omega A}{c} = \frac{\Omega L^2}{\sqrt{3} c}$$

The condition to sustain the laser oscillation (given in the problem),

$$v_{\pm} = \frac{m}{L_{\pm}} c, \text{ } m = 1, 2, 3, \dots \text{ integers} \quad (1 \text{ point})$$

$$\Delta v = v_{-} - v_{+} = \frac{m}{L_{-}} c - \frac{m}{L_{+}} c \approx mc \frac{\Delta L}{L^2} = v \frac{\Delta L}{L} \quad (1 \text{ point})$$

the approximation arises from $L_+L_- \approx L^2$

where L is the perimeter of the triangular ring. Hence,

$$\Delta v = \frac{\Delta L}{L} v = \frac{4A}{Lc} v \Omega = \frac{1}{\sqrt{3}} \frac{L}{\lambda} \Omega \quad (1 \text{ point})$$

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III. Plasma Lens

- a) Consider cylindrical Gaussian surface of radius r and length ℓ about the central axis.

From Gauss's law

$$E 2\pi r \ell = \frac{1}{\epsilon_0} (\text{charge inside}) \quad (0.5 \text{ points})$$

$$= -\frac{1}{\epsilon_0} ne(\pi r^2 \ell)$$

$$E_r = -\frac{ner}{2\epsilon_0} \quad (0.5 \text{ point})$$

- b) From Ampere's law, $B_\theta = -\frac{\mu_0 n e r v}{2}$ (2 points)

- c) The net Lorentz force is

$$\vec{F} = \left(\frac{ne^2 r}{2\epsilon_0} - \frac{\mu_0 ne^2 r v^2}{2} \right) \hat{r} = \frac{ne^2 r}{2\epsilon_0} \left(1 - \frac{v^2}{c^2} \right) \hat{r}$$

where $c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$ (1 point)

- d) $F_r \rightarrow 0$ as $v \rightarrow c$, this implies the electric force and magnetic force cancel each other out. (1 point)

- e) The stationary plasma particles have $v = 0$, hence $F_{r'} = \pm eE_r$.

where $E_{r'} = -\frac{neR^2}{2\epsilon_0 r'} + \frac{n_o e r'}{2\epsilon_0}$

for positive ion $F_{r'} = -\frac{ne^2 R^2}{2\epsilon_0 r'} + \frac{n_o e^2 r'}{2\epsilon_0}$

for electron $F_{r'} = \frac{ne^2 R^2}{2\epsilon_0 r'} - \frac{n_o e^2 r'}{2\epsilon_0}$

and there is no cancellation from the magnetic force.

As a result the plasma electrons will be blown out, and the ions are pulled in. (2 points)

- f) The net force on the electron beam in plasma medium is given by,

$$\vec{F} = \frac{ne^2 r}{2\epsilon_0} \left(1 - \frac{v^2}{c^2} \right) \hat{r} - \frac{n_o e^2 r}{2\epsilon_0} \hat{r} \quad (2 \text{ points})$$

in the limit $v \rightarrow c$, $\vec{F} \approx -\frac{n_o e^2 r}{2\epsilon_0} \hat{r}$ (1 point)